

Crystal momentum $\hbar\vec{k}$

$$i\hbar \frac{d}{dt} \langle A \rangle = \langle AH - HA \rangle$$

translation operator $T\psi(x) = \psi(x+a)$

$\langle T \rangle$ is a constant of motion for a perfect crystal

$$i\hbar \frac{d}{dt} \langle T \rangle = \langle TH_0 - H_0T \rangle = 0$$

Consider an external force in the x-direction

$$F_{ext} = -\frac{dU}{dx} \Rightarrow U = -F_{ext}x$$

$$i\hbar \frac{d}{dt} \langle T \rangle = \langle T(H_0 + U) - (H_0 + U)T \rangle$$

$$i\hbar \frac{d}{dt} \langle T \rangle = \langle T(H_0 - F_{ext}x) - (H_0 - F_{ext}x)T \rangle$$

$$i\hbar \frac{d}{dt} \langle T \rangle = \langle -TF_{ext}x + F_{ext}xT \rangle = \langle -F_{ext}(x+a)T + F_{ext}xT \rangle$$

Crystal momentum

$$i\hbar \frac{d}{dt} \langle T \rangle = \langle -F_{ext} (x+a)T + F_{ext} xT \rangle$$

$$i\hbar \frac{d}{dt} \langle T \rangle = \langle -F_{ext} aT \rangle = -F_{ext} a \langle T \rangle$$

$$\langle T \rangle = \langle e^{-ikx} u_k(x) | T | e^{ikx} u_k(x) \rangle = \langle e^{-ikx} u_k(x) | e^{ik(x+a)} u_k(x+a) \rangle$$

$$\langle T \rangle = e^{ika} \langle e^{-ikx} u_k(x) | e^{ikx} u_k(x) \rangle = e^{ika}$$

$$i\hbar \frac{d}{dt} e^{ika} = -F_{ext} a e^{ika}$$

$$i\hbar(ia) \frac{dk}{dt} e^{ika} = -F_{ext} a e^{ika}$$

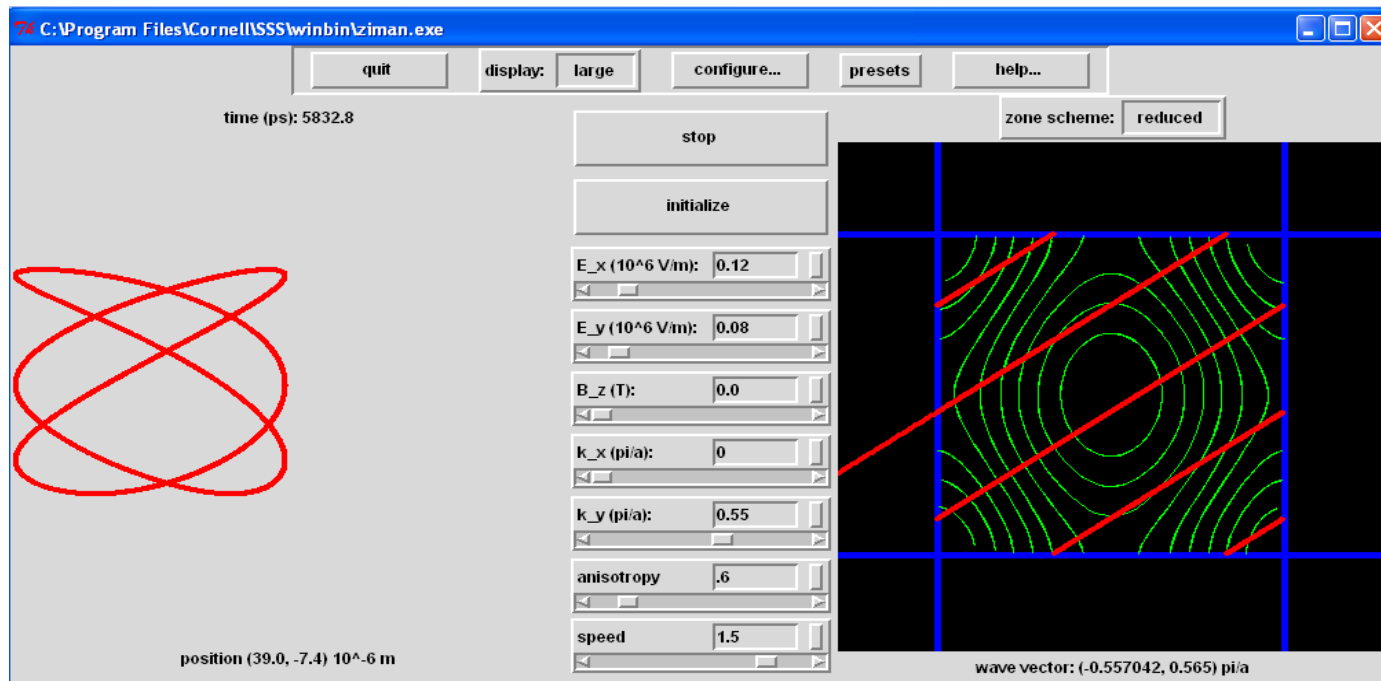
$$\hbar \frac{dk}{dt} = F_{ext}$$

Crystal momentum

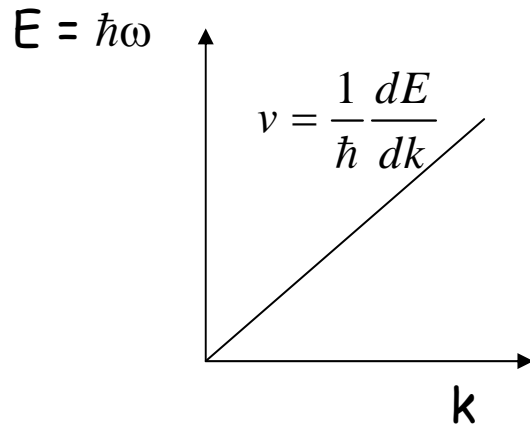
$$\hbar \frac{d\vec{k}}{dt} = \vec{F}_{ext}$$

$$\frac{d\vec{p}}{dt} = \vec{F}_{total}$$

$$\hbar\vec{k} \neq \vec{p}$$



Group velocity

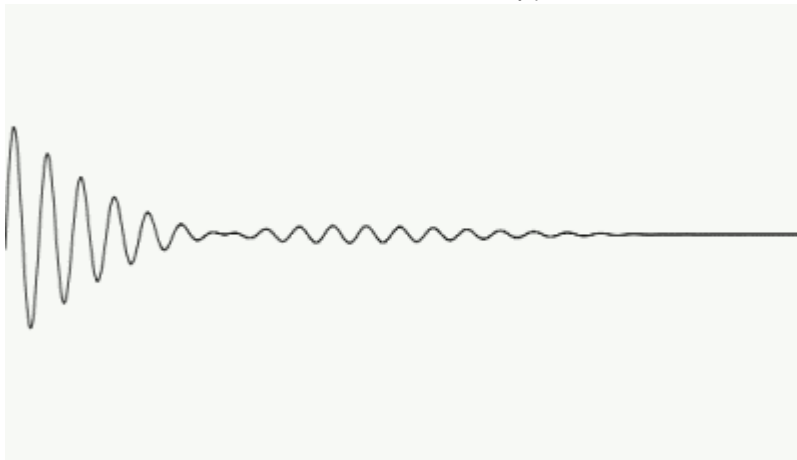


$$v_g = \frac{1}{\hbar} \frac{dE}{dk}$$

$$\frac{dv_g}{dt} = \frac{1}{\hbar} \frac{d^2E}{dk dt} = \frac{1}{\hbar} \frac{d^2E}{dk^2} \frac{dk}{dt}$$

$$\frac{dv_g}{dt} = \frac{1}{\hbar^2} \frac{d^2E}{dk^2} F_{ext}$$

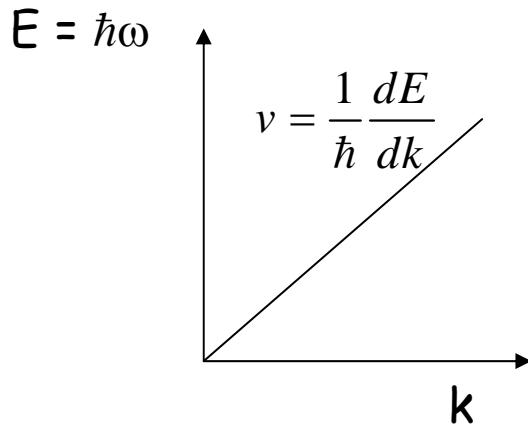
$$F_{ext} = \hbar^2 \left(\frac{d^2E}{dk^2} \right)^{-1} \frac{dv_g}{dt} = m^* a_g$$



v_g is the velocity of a wave packet
 $\Delta x \Delta k \sim 1$

Particles in a semiconductor can be thought of as free particles with an effective mass.

Group velocity

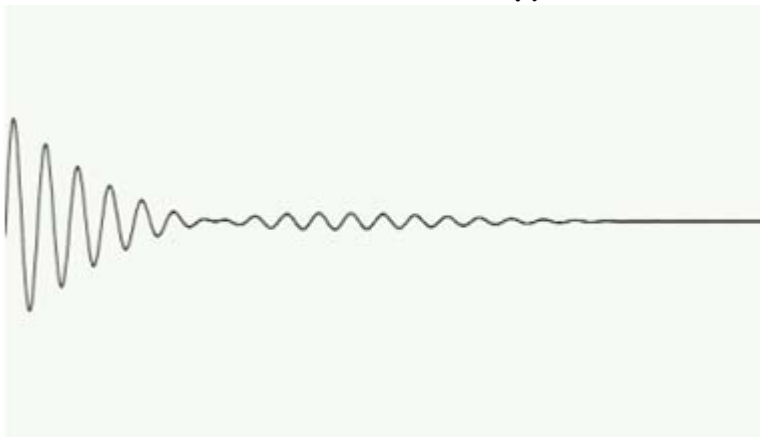


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$$a_g = \frac{dv_g}{dt} = \frac{1}{\hbar^2} \frac{d^2E}{dk^2} F_{ext}$$

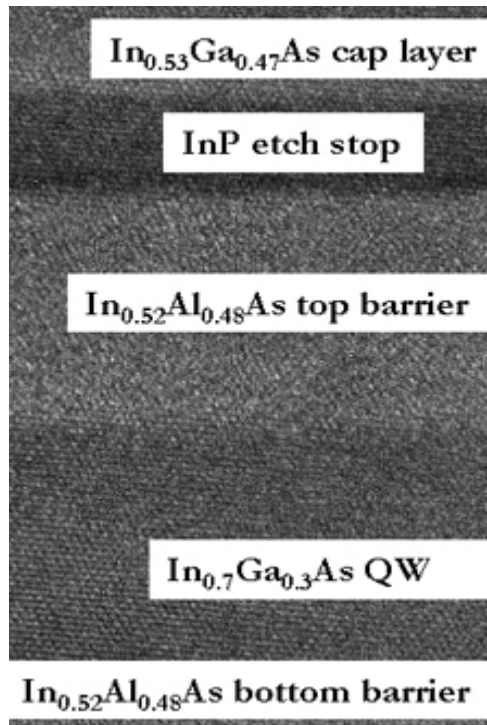
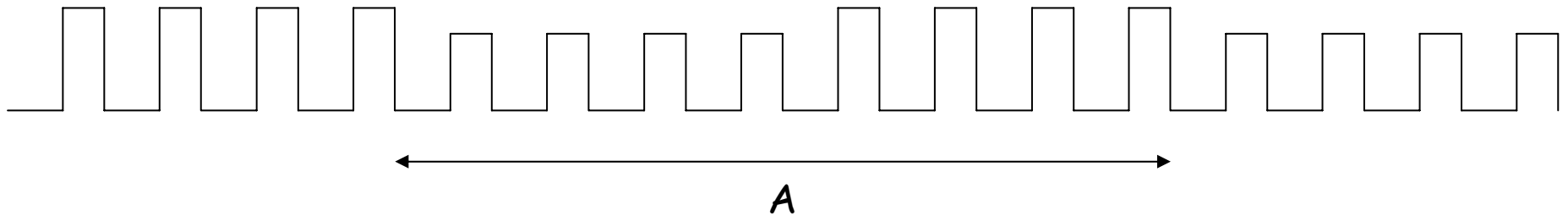
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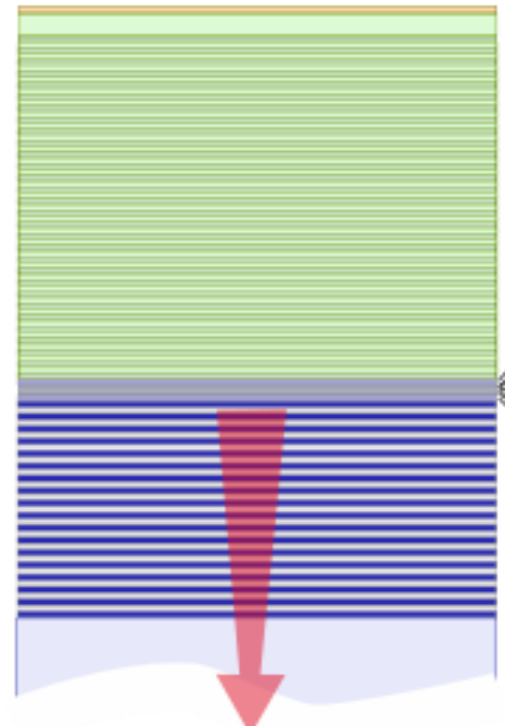
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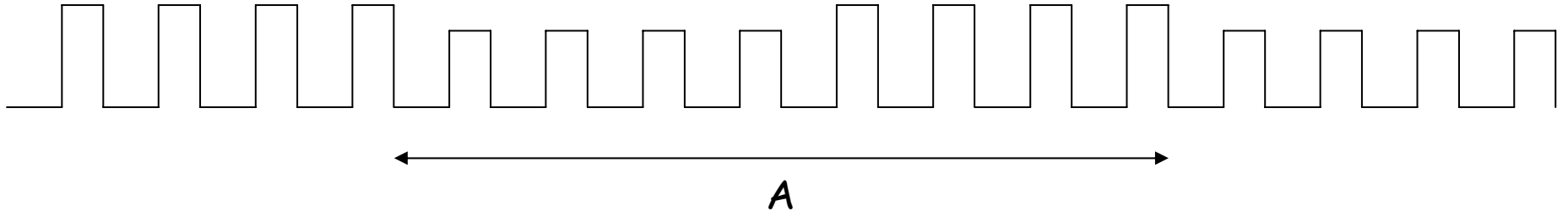
Mini bands in superlattices



laser →



Bloch oscillations



$$\varepsilon(k) = -\varepsilon_0 \cos(kA)$$

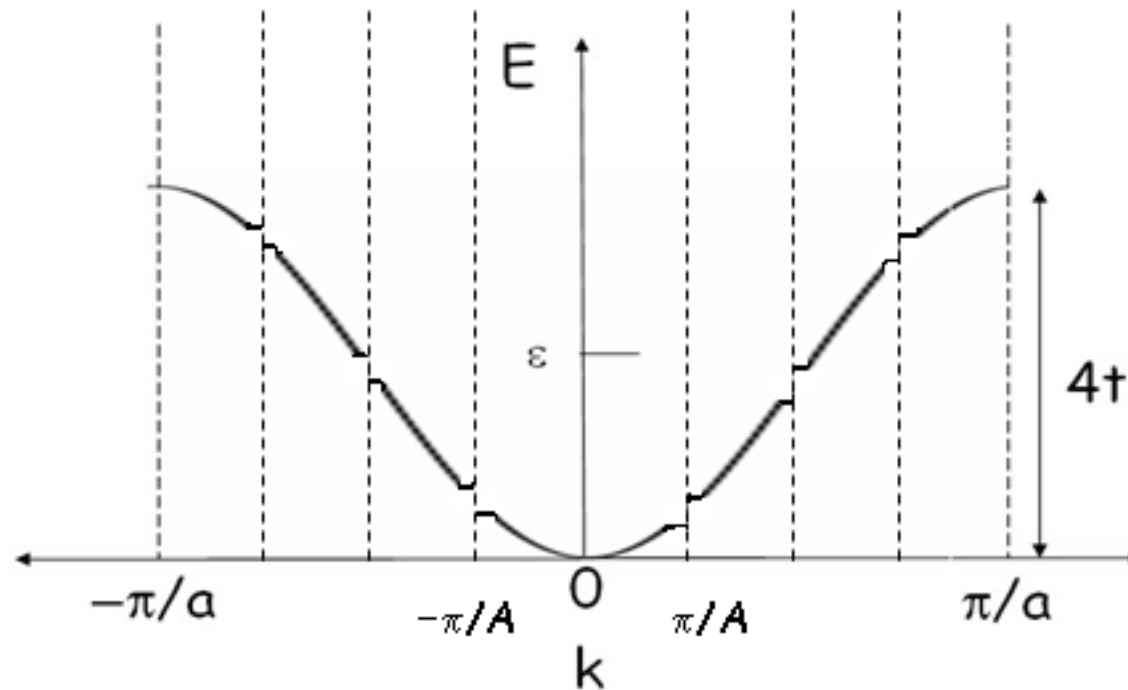
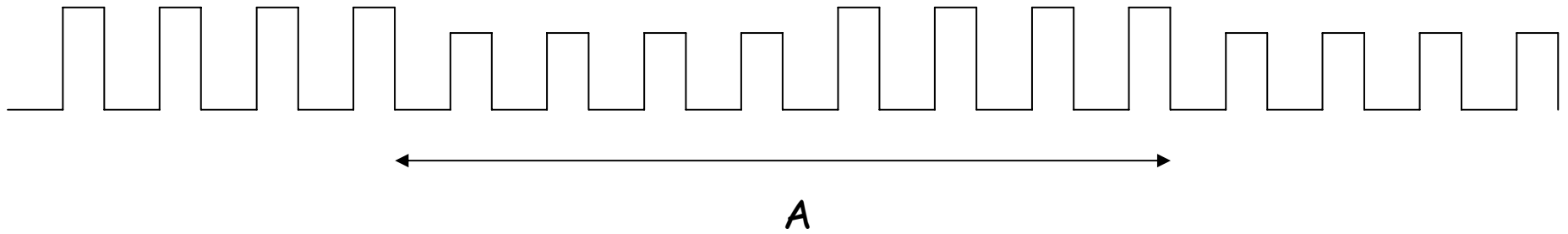
$$v(k) = \frac{1}{\hbar} \frac{d\varepsilon}{dk} = \frac{A\varepsilon_0}{\hbar} \sin(kA)$$

$$x = \int v dt = \int v(k) \frac{dt}{dk} dk \quad -eE = \hbar \frac{dk}{dt} \quad \Rightarrow \quad k = -\frac{eEt}{\hbar}$$

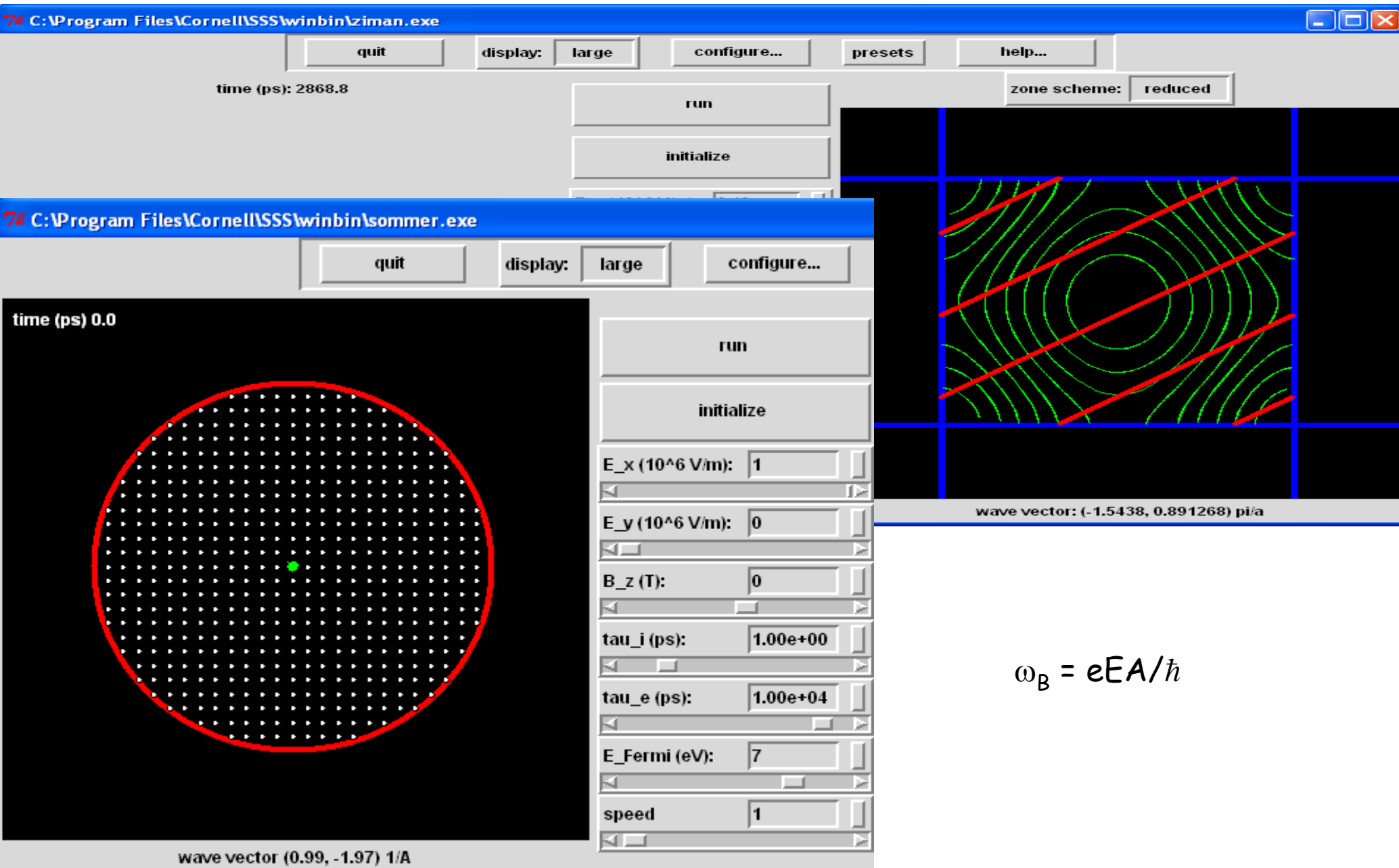
$$x = \int \frac{-\hbar}{eE} \frac{A\varepsilon_0}{\hbar} \sin(kA) dk = \frac{\varepsilon_0}{eE} \cos(kA) = \frac{\varepsilon_0}{eE} \cos\left(\frac{eEA t}{\hbar}\right)$$

The Brilluoin zone is much smaller for a superlattice. This makes the Bloch oscillation frequency $\omega_B = eEA/\hbar$ higher.

Mini bands in superlattices



Bloch oscillations



Thermoelectric effects

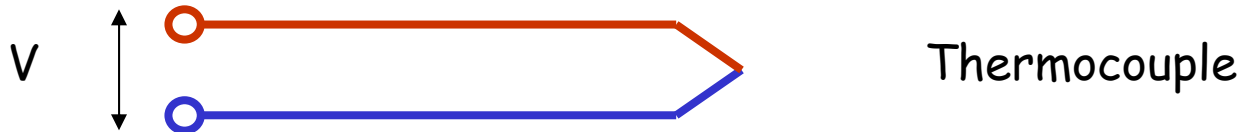
Seebeck effect:

A thermal gradient causes a thermal current to flow. This results in a voltage which sends the low entropy charge carriers back to the hot end.

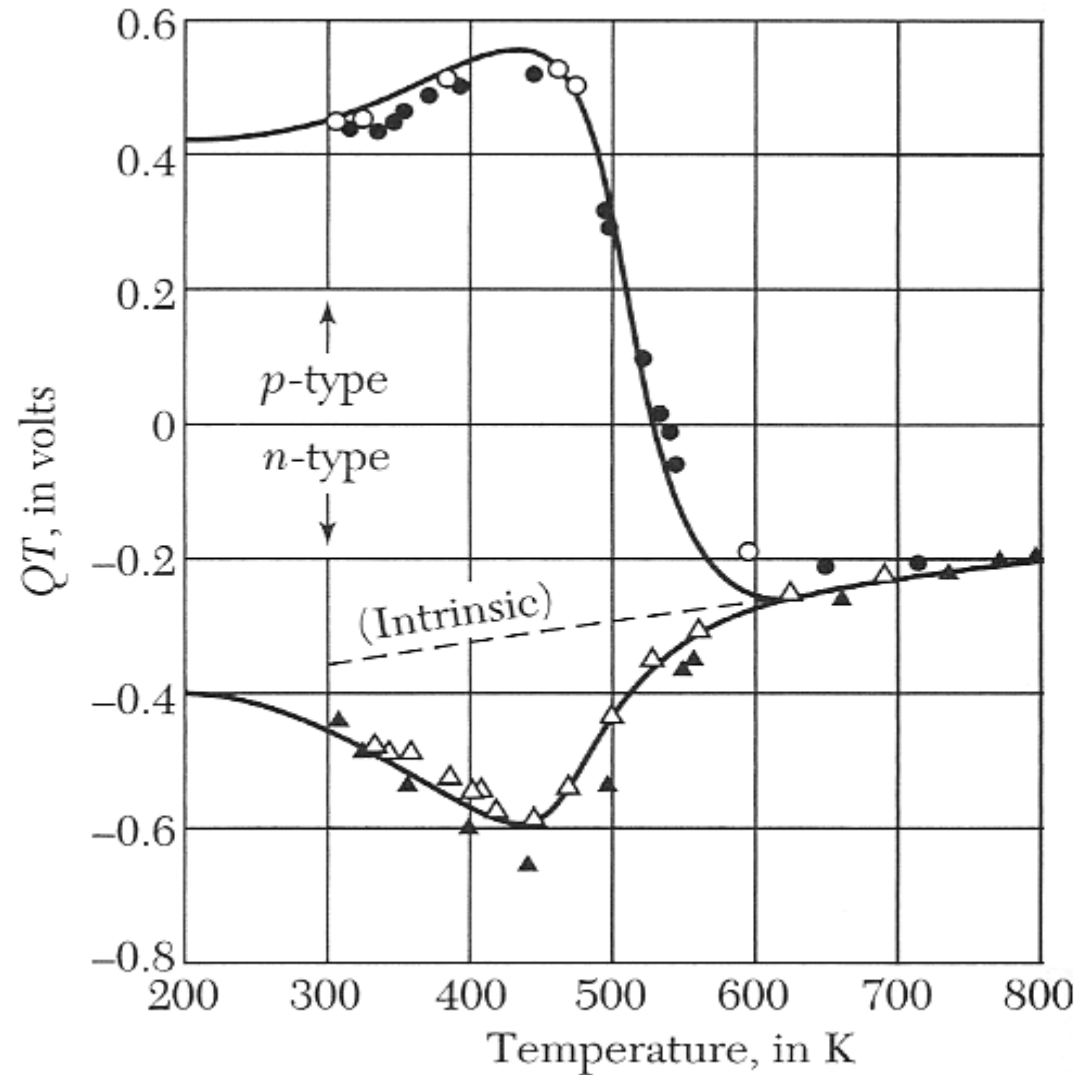
$$\vec{E} = Q\nabla T$$

Q is the absolute thermal power. The sign of the voltage is the same as the sign of the charge carriers.

The Seebeck effect can be used to make a thermometer.



Thermoelectric effects



Thermoelectric effects

Electrons carry charge, energy, and entropy. The charge is $-e$. The entropy has to be calculated using statistical mechanics.

The internal energy per electron can be calculated using:

$$u = \frac{U}{N} = \int_{E_c}^{\infty} ED(E)f(E)dE$$

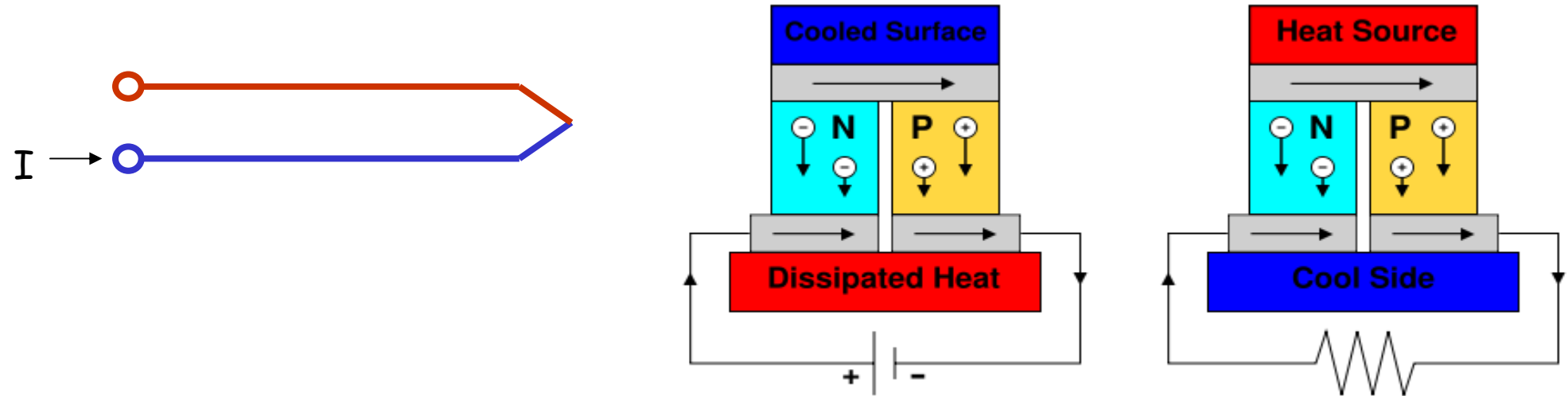
The energy current is related to the electrical current by the Peltier coefficient.

$$j_U = \Pi j_q$$

Π is the Peltier coefficient.

Thermoelectric effects

Peltier effect: driving a current through a bimetallic junction causes heating or cooling.



Cooling takes place when the electrons make a transition from low entropy to high entropy at the junction.